

Introducing remark

In early days of QM intense discussion of philosophical issues (the analysis of concepts, principles and ultimate presuppositions) of the theory. This culminated in 1935 with EPR debate between Einstein and Bohr. Bohr (and Copenhagen interpretation) were officially declared the winners. Bohr reigned supreme until the 1960s, except for the efforts of Bohm to construct h.v. models of QM. This was partly due to the influence of von Neumann's no-go theorem proved in the oft-cited but little read *Mathematical Grundlagen der Quantenmechanik* (1932).

Main technical advance was Gleason's theorem (1957) on possible prob. measures that could be defined on its projection lattice of a Hilbert space. This Gordon proved to be as they to the work of Bell and Finkelstein and Speiser (1967). A parallel development was the discovery of a non-locality property of h.v. states in 1964. This was the Bell inequalities which led to many experimental tests culminating in the work of Aspect, Dalibard and Roger about 1980 's onward.

Notice different reactions here as between Physicists & Philosophers.

Aspect's result expected by physicists - regarded by many philosophers as revolutionary
Minimal instrumentalist (pragmatic) interpretation of QM in
 used by most physicists. Philosophers want to understand

N.B. $2P_{\text{system}}$ is a N.D. wave function, but it is the exact fun of the relative norm of the 2-body spin wave function relative to an appropriate phase of reference for a slowly moving 2-body system

In order to proceed we distinguish 3 versions of the individual particle interpretation of QM wheel call A, B & C.

Thus given

3 means an observable not in an eigenstate

- A) Sharp unknown value
- B) unsharp or 'fuzzy' value
- C) Undefined or meaningless value.

QM is not mysterious, mysterious with new concepts like potentiality or latency, and mysterious in recognizing limitations on the applicability of old concepts.

3 means a latent measurement does

- A) Pre-existing value revealed
- B) Potential & latent value actualized
- C) Undefined value becomes defined.

How is this change brought about?

- 1) Interaction with a physical instrument
- 2) Action of human consciousness

The EPR Argument (1935)

This is an argument for the incompleteness of QM.

EPR takes as necessary condition for

Completeness: Every element of reality has a counterpart in the theory, associated with it.

They take as sufficient condition for an element of reality: If without disturbing a system we can predict with certainty the value of an observable already, before the prediction, then exists an element of reality corresponding to that observable.

The EPR argument is based operationally on wave B and shows that this wave is incomplete in the sense that even in non-simultaneous observations do have sharp values (not referred to in pointer B).

So the argument writes itself in terms of pointer A.

But the argument involves a locality assumption

Locality: Unsharp \rightarrow sharp at a distance

(Disturbance Bell and Further zone of disturbance)

Bohr's response to EPR was to query whether we should not well insist upon C and use
locality₂: undefined \rightarrow defined
 'at-a-distance' forbidden

But this does purport to deny that no physical non-locality involved.

EPR showed that
 $QM + \text{locality}_2 \rightarrow$ Incompleteness
 (assuming near B)

a QM \rightarrow Non-locality, or Incompleteness

Yes is the Einstein Dilemma ^{posed by EPR}
 But Bell showed (1964) $\left\{ \begin{array}{l} \text{EPR} \text{ leads to } \text{EPR Paradox} \\ \text{The } \text{EPR} \text{ leads to } \text{EPR Paradox} \end{array} \right.$

Completion \rightarrow non-locality₃
 (i.e. near A)

where.

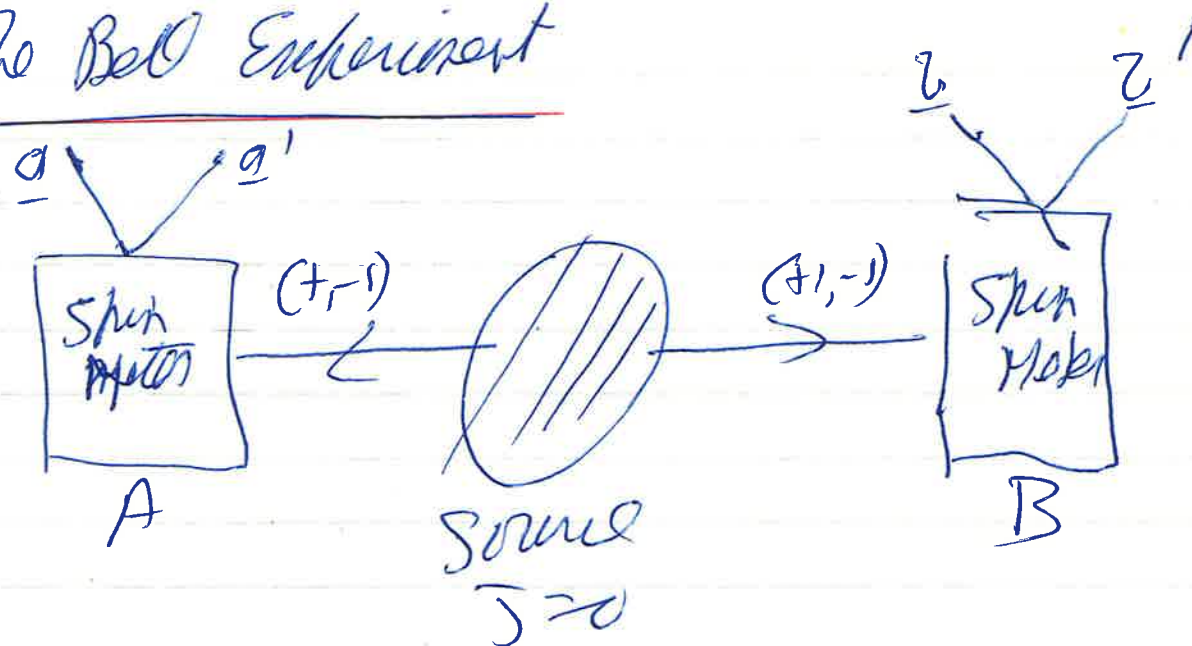
locality₃ sharp \rightarrow sharp
 at-a-distance forbidden

The Bell Argument

locality₃ \rightarrow Bell Inequality
 which is contradicted by
 QM (and by experiment)

cf a particular Aspect, Dalibard and
Roger (P.R.L. (1982))

The Bell Experiment



a_n denotes spin-component of A-particle parallel to direction \underline{a} for n^{th} pair of particles
 similar a_n', b_n, b_n'

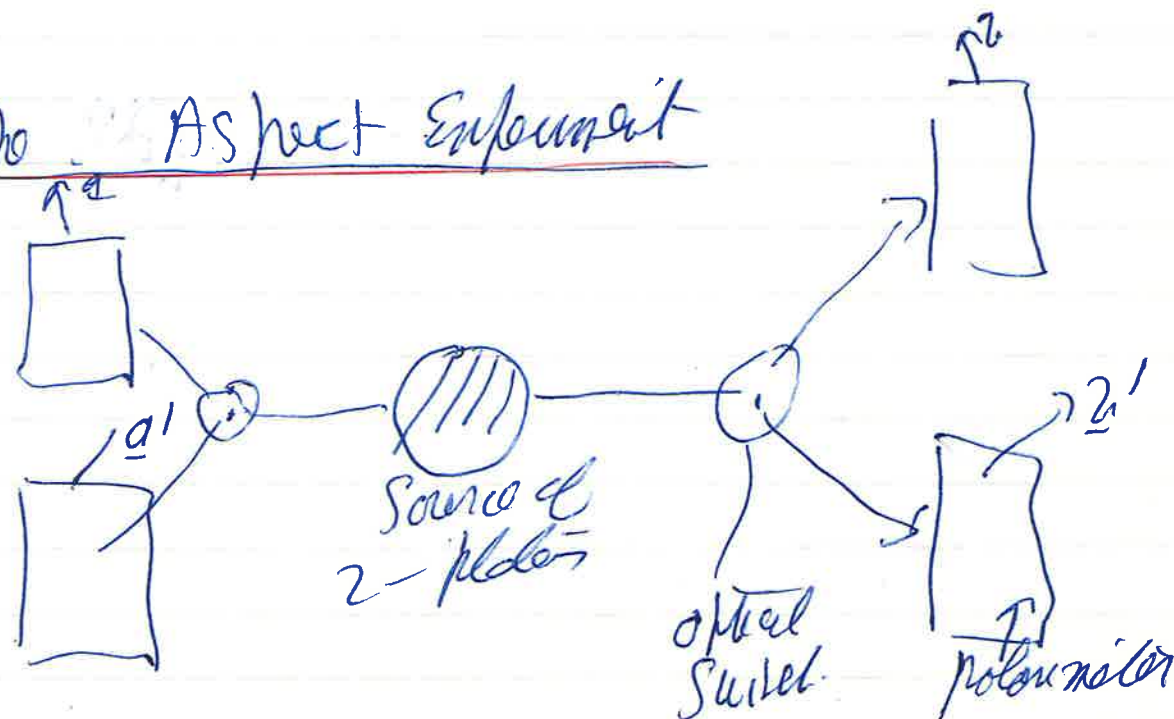
then define correlation coefficients

$$C(a, b) = \frac{1}{N} \sum_n a_n b_n, \text{ etc}$$

The Bell Inequality reads

$$|C(a, b) + C(a, b') + C(a', b) - C(a', b')| \leq 2.$$

The Aspect Experiment



The Stapp-Chorhard Approach

Can we reformulate argument for Bell Inequality in terms of measurement results so it would apply also to user B.

Locality₄: Classical state of macroscopic object cannot be altered at a distance.

Principle of Local Counterfactual Definiteness (PLCD)

Result of an unperformed experiment has a definite result which does not depend on the setting of a remote piece of apparatus.

(So $\sim \text{PLCD} \rightarrow \sim \text{Loc}_4$ for Stapp & others)

But in a genuinely indeterministic situation (as in user B) PLCD is suspect.

Contrast a) clock striking as I raise my hand
b) pen dropping as I raise my hand.

Even if some experiment in a new world report same result (counterfactual).

Statistical Nonlocality

No statistical effects produced at a distance - tagging information at the wrong location - Oxford lecture, notes in Chelsea example.

- Impossibility of the Bell Telephones

Is locality violated in EPR?

Violates	loc ₁	loc ₂	loc ₃	loc ₄	loc ₅
view					
A	no	no	yes	yes	no
B	yes	no	no	no	no
C	no	yes	no	no	no.

N.B. Locality problem does not arise in statistical example
~~offered since loc₅ not violated~~

Query Does Violation of Locality conflict with S.R.?

Violation of loc₃ now remains for loc₁
 cf Shimony (1978) conceptual consistency in case of violations of loc₁.

Does S.R. prohibit instantaneous effects
 Tachyon theory, Causal loop paradox

Kochen-Specker Paradox

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So with zero from of non-locality and
can retain realism. But can we?
There is another difficulty posed for
realism by the Kochen-Specker paradox.
They demonstrated a purely algebraic
problem about attributing values to
observables.

Kochen & Specker showed the following.

view $A + (A^2 + (FVVO)) \Rightarrow \text{contradiction.}$

[VR says. $P_\phi(\lambda) = 0 \Rightarrow [\phi]^\phi \neq \lambda.]$
(Sum of 2, FM.)

FUNE says. If \hat{A} and \hat{B} are two Hermitian
and there exists a function f s.t.
 $\hat{A} = f(\hat{B})$ then

$$[\hat{A}]^\phi = f([\hat{B}]^\phi)$$

$$10. [f(\hat{B})]^\phi = f([\hat{B}]^\phi)$$

Proof relies on FUNE holding for

$$\hat{A} = f(\hat{B}), \quad \hat{A} = g(\hat{C}) \quad [\hat{B}, \hat{C}] \neq 0.$$

$$\text{So } [\hat{A}]^\phi = f([\hat{B}]^\phi)$$

$$= g([\hat{C}]^\phi)$$

— constraint
on values for
incompatible magnitudes

N.B. If \hat{B} & \hat{C} are measured, \hat{A} must
be degenerate for the situation to
arise.

Insertion
*

Extension of notation to two separated systems:

$$[\mathcal{Q} \times \mathcal{I}]_{\{\langle A, B \rangle\}}^{\Phi} (\mathcal{D}, \mathcal{E})$$

We can break FUNC and retain realism by supposing that each nat. speaks two languages more than one may speak

Then we distinguish A_B, A_C

$$\text{Here } [A_B]_{\{B\}}^{\phi} \stackrel{\text{iff}}{=} H(\{B\})$$

$$[A_C]_{\{C\}}^{\phi} \stackrel{\text{iff}}{=} G(\{C\})$$

In general we shall assume FUNC*
which says: effectives

$$A_B = A_C, \text{ if } \vec{B}, \vec{C} \text{ are commensurable}$$

measured by the same
i.e. $\exists g \text{ s.t. } \vec{C} = g(\vec{B})$
 $g \in 1:1$

Depends on $\{B\}$ equivalent class of
all C s.t. $\exists g \text{ s.t. } C = g(B)$

$$\text{Then we write } [A]_{\{B\}}^{\phi} \stackrel{\text{iff}}{=} [A_B]_{\{B\}}^{\phi}$$

N.B. FUNC* is consistent with and
independent of VR.

This replicates intentional Contextuality

But the value $[A]_{\{B\}}^{\phi}$ may also
depend on the environment, in particular
the affector set to measure some
magnitudes C

$$\text{So we write } [A]_{\{B\}}^{\phi}(C)$$

Then replicates Environmentally Contextuality

* insert

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We now apply these notions to
two separated systems.
For all Q, A, B, C, D, E , where $Q = h(A)$ and
 A, B, C, D and E are all nominal

OLOC

$$[Q \otimes I]_{\{A, B\}}^{\phi} (D, E) = [Q \otimes I]_{\{A, C\}}^{\phi} (D, E)$$

ELOC

$$[Q \otimes I]_{\{A, B\}}^{\phi} (D, E) = [Q \otimes I]_{\{A, B\}}^{\phi} (D, E)$$

where $\langle A, B \rangle$ is the nominal physical
magnitudes of the first system
where associated operator \hat{O} is constructed
by $\hat{A} \otimes \hat{B}$ by the relation

$$\hat{O} = \sum_{i,j} c_{ij} \hat{P}_i \otimes \hat{P}_j$$

where $\hat{A} = \sum \alpha_i \hat{P}_i$ $\hat{B} = \sum \beta_i \hat{P}_i$
and $c_{ij} = F(\alpha_i, \beta_j)$ where $F: \mathbb{R}^2 \rightarrow \mathbb{R}$ (1:1)

In words-

OLOC locally nominal physical magnitudes
on either of two spatially separated
systems are not split, they
are logically contextually related to
the specification of different nominal
physical magnitudes for the first system

FLOC : the value possessed by a local physical magnitude cannot be changed by altering the arrangement of a remote piece of apparatus which forms part of its measurement context for the concerned system.

VR ① we do not require LOC in specifying FLOC but FLOC is as a probability of locally particles of LOC states.

② In terms of measurement results $LOC \equiv FLOC$ but demonstrate a dependence of outcomes recorded by apparatus connected to one system on the setting of the apparatus connected to the other (remote) system.

Redhead - Wesleywood Thesis

We shall show that

* $FLOC \rightarrow VR \rightarrow FLOC \rightarrow LOC$

Define Value Rule VR $\Rightarrow P_i^{\Phi}(\lambda) = 0 \Rightarrow \{ \Phi \neq \lambda \}$. (OK-S type)
Preliminary Comments

Proofs of nonlocality via Bell Inequality have been charged with 'hidden' assumptions. AS Fine remarked (1974) 'Hidden variables are no thing, hidden assumptions another'.

Fine has claimed (1974) that Bell's
 original type of proof involved FINE
 in the form of the Product Rule (JD
 assumption that QM joints and phase-space
 joints were identical)
 However Fine said we are committed
 to joint distributions for non-commuting
 observables even in the Eberhard type
 of proof (1982) (in h.v. approach this is true)
 But Fine is mistaken here.

Bell \rightarrow KQM-joints
 KQM-joints \rightarrow self-joints in a model.
 But this model need not be the real
 world!

However all proof of non-locality via
 Bell's theorem do involve notions
 of probabilities theory that might be
 challenged.

Question Can we give a purely algebraic
 proof of non-locality?

Question posed by Bub (1976) in the form:

Is it possible to extend Maczynski's
 Theorem (1971) to locally maximal
 observables?

M's Theorem says we do not need
 ontological contextualism to allow
 value assignments to maximal observables

If M's theorem cannot be extended then
 would mean
 locally maximal observables must be
 ontologically contextual \Rightarrow validation of QM

Unfortunately Demopoulos, Bub, & Humphreys
have shown that H's theorem -
can be extended to locally maximal
assemblies.

So OLOC need not be violated

Our next show however that if
we retain Funca, VR then
OLOC & ELOC are both most fail.

Violating OLOC means we cannot
specify a locally maximal observable
independently of properties relating to
the whole combined system
— leads to alleged holism in which
it is impossible to make sense of
a realist construal of QM that considers
properties independently with each of
two separated systems

Violating ELOC (if OLOC is assumed)
is the fact of nonlocality envisaged
in the Bell-type nonlocal arguments.

The steps in the proof of * are as follows:

- 1.) Derive EVR from VR and Func^*
 EVR of \hat{Q}_1 and \hat{Q}_2 commute and \hat{P} is
 maximal operator s.t. $\hat{Q}_1 = f(\hat{P})$, $\hat{Q}_2 = g(\hat{P})$
 Then $P_{\hat{Q}_1, \hat{Q}_2}^\Phi(\lambda, \mu) = 0 \Rightarrow$

$$[\hat{Q}_1]_{\{\hat{P}\}}^\Phi(R) \neq \lambda$$

$$\vee [\hat{Q}_2]_{\{\hat{P}\}}^\Phi(R) \neq \mu.$$

N.B. EVR for general commuting \hat{Q}_1, \hat{Q}_2
 is not derivable from VR and Func^*
 act. $\Rightarrow \text{FUNC}$ is contradictory
 This was proved by Foul on 1974.

- 2.) As a special case of EVR we
 obtain, if $\hat{Q} = h(\hat{A})$, $\hat{Q}' = h'(\hat{B})$
 where \hat{A}, \hat{B} are maximal
 $P_{\hat{Q} \otimes I, I \otimes \hat{Q}'}^\Phi(\lambda, \mu) = 0$
 $\Rightarrow [\hat{Q} \otimes I]_{\{\langle \hat{A}, \hat{B} \rangle\}}^\Phi(A, B) \neq \lambda$
 $\vee [I \otimes \hat{Q}']_{\{\langle \hat{A}, \hat{B} \rangle\}}^\Phi(A, B) \neq \mu$

3.) Incompatibility of EVR and Locality

Consider correlated state of 2 separated
 systems

$$|\Psi\rangle = \sum C_m |a_m\rangle \otimes |b_m\rangle$$

\hat{A} has eigenvalues a_1, \dots, a_N
 \hat{B} has eigenvalues b_1, \dots, b_N

Consider nonmaximal operator \hat{Q} s.t.
 $\hat{Q} = f(\hat{A}) = g(\hat{A}')$ for maximal \hat{A}, \hat{A}'
 where $[\hat{A}, \hat{A}'] \neq 0$.

Then it is easy to show that

$$[A \otimes I]_{\{\langle A, B \rangle\}}^4 (A, B) = a_m \quad (1)$$

$$\Rightarrow [I \otimes B]_{\{\langle A, B \rangle\}}^4 (A, B) = b_m.$$

and $[I \otimes B]_{\{\langle A', B \rangle\}}^4 (A', B) = b_m$

$$\Rightarrow [f(A \otimes I)]_{\{\langle A', B \rangle\}}^4 (A', B) = f(a_m) \quad (2)$$

Now apply ELOC & OLOC to eqn (2)

$$[I \otimes B]_{\{\langle A, B \rangle\}}^4 (A, B) = [I \otimes B]_{\{\langle A', B \rangle\}}^4 (A', B)$$

where we obtain FUNC^{xx}

$$[f(A \otimes I)]_{\{\langle A', B \rangle\}}^4 (A', B) = f([A \otimes I]_{\{\langle A, B \rangle\}}^4 (A, B))$$

This is not quite FUNC which would have same environmental context on both sides of \cdot .
 But can we do this to demonstrate a K-S type contradiction by following this argument parallel to that used by K-S for smallest state of 2 spin-1/2 particles in their 1967 paper.

Comments

- 1) Fine (1977) and Stairs (1978) claimed CVR was contradictory

But we have derived CVR on a theorem from VR and Fare^x.

Explanation CVR would be contradictory if we did not allow for violation of locality - this is shown in all work of Fine and Stairs

- 2) Minimal, transparent use of probability theory. On user A QM is non-local

- 3) What about Stochastic h.v. theories where measurement results are stochastic related to the hidden state of the microsystem (so FM is given up)? The Eberhard-type proof now fails because of PLED being inapplicable. VR is not applicable to our approach does not apply.

Bell Theorem can no longer give us identity locality = factorizability.

Perhaps we just have to give up factorizability - cp Fine.

Final remark

Interaction of Physics and Philosophy

- either 1) Physics should be changed in the light of philosophical precepts
or 2) Philosophical ideas should be modified in the light of new theories in physics.

Role of Philosophy of Physics as I see it is to show what total packages are available — when choice is matter for argument but not decisively settleable.

(I deal with on Physics) — side (c)

The role of philosophy in physics then is to promote understanding of what physics really is committing us to.

Of course theoretical physicists do use philosophical analysis of a sort in their work — Einstein and Bohr are notable examples. However in the case of Einstein general relativity notoriously fails to implement his philosophy, while Bohr gives us a very weaker mesh — mesh of Kierkegaard, Kant and William James, the mottled zone of his former.

But the surprising thing is that muddled philosophy is quite consistent with splendid physics.

In a sense one could claim that Einstein did not understand Relativity

not Bohr quantum mechanics, but
then one can do physics without
understanding it - rather like riding
a bicycle without knowing anything
about rigid-body dynamics.
But how much better to do
physics and understand it - to
combine the study of physics
with the study of philosophy of
physics
